Formulation and Solution Strategies

Outline

- Review of Field Equations
- Types of BCs
- BCs on Coordinate Surfaces
- BCs on Oblique Surface
- Line of Symmetry BCs
- Interface BCs
- Problem Classification
- Stress Formulation
- Displacement Formulation
- Principle of Superposition
- Uniqueness of Elastic Solution
- Saint-Venant's Principle
- Solution Strategies
- Mathematical Techniques

Governing Equations in linear elasticity

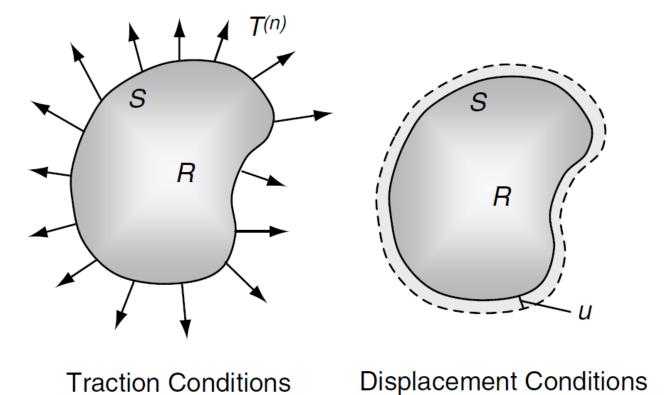
- Strain-displacement relations: $\varepsilon_{ij} = \frac{1}{2} (u_{i,j} + u_{j,i})$ (6 eqns)
- Strain compatibility: $\varepsilon_{ij,kl} + \varepsilon_{kl,ij} \varepsilon_{ik,jl} \varepsilon_{jl,ik} = 0$ (6 eqns)
- Equilibrium: $\sigma_{ii,j} + F_i = 0$ (3 eqns)
- Isotropic Hooke's Law:

$$\sigma_{ij} = \lambda \varepsilon_{kk} \delta_{ij} + 2G \varepsilon_{ij}; \qquad \varepsilon_{ij} = \frac{1+\nu}{E} \sigma_{ij} - \frac{\nu}{E} \sigma_{kk} \delta_{ij}. \qquad (6 \text{ eqns})$$

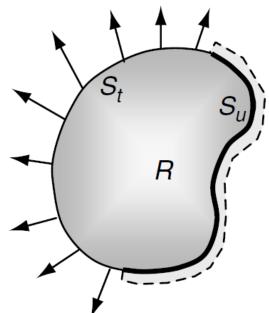
- 15 equations for 15 unknowns (3 displacements, 6 strains, 6 stresses).
- May define the entire system as

$$\int \{u_i, \varepsilon_{ij}, \sigma_{ij}; \lambda, G, F_i\} = 0$$

Traction and Displacement Boundary Conditions



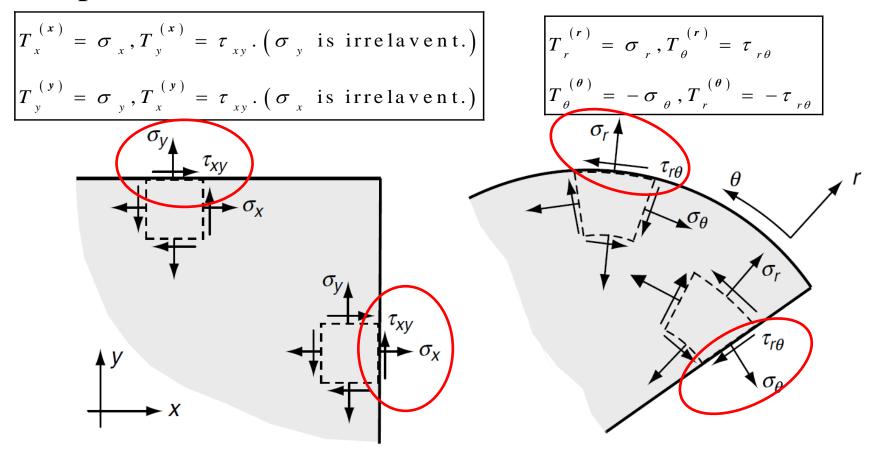




Mixed Conditions (c)

Boundary Conditions on Coordinate Surfaces

• The traction specification can be reduced to a stress specification.

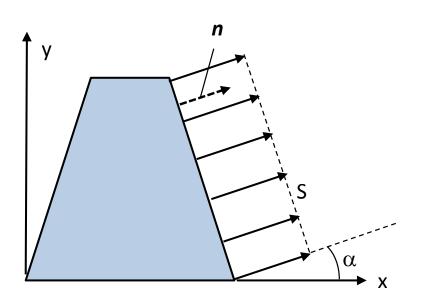


(Cartesian Coordinate Boundaries)

(Polar Coordinate Boundaries)

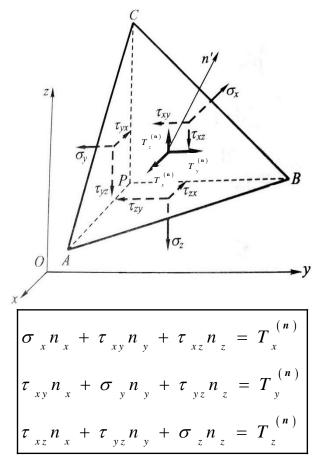
Boundary Conditions on Oblique Surfaces

• On general non-coordinate surfaces traction vector will not reduce to individual stress components and general form of traction vector must then be used.



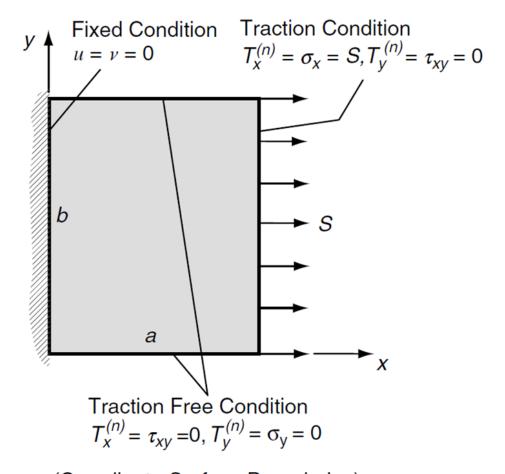
$$\sigma_{x} n_{x} + \tau_{xy} n_{y} = T_{x}^{(n)} = S \cos \alpha$$

$$\tau_{xy} n_{x} + \sigma_{y} n_{y} = T_{y}^{(n)} = S \sin \alpha$$



Example of Boundary Conditions

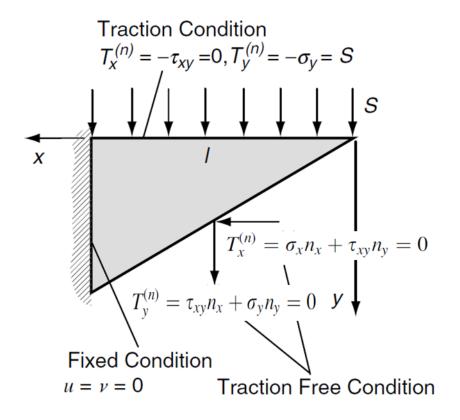
Example (1)



(Coordinate Surface Boundaries)

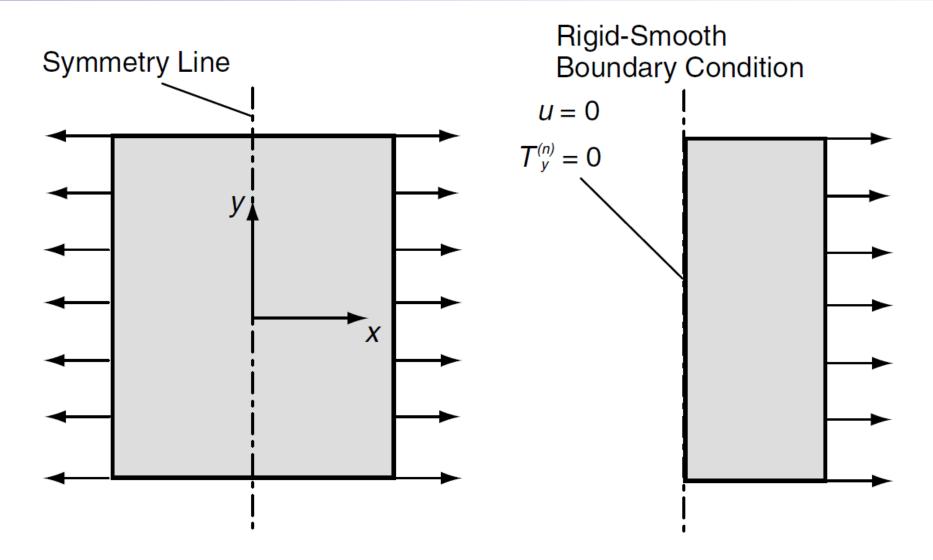
Example of Boundary Conditions

Example (2)



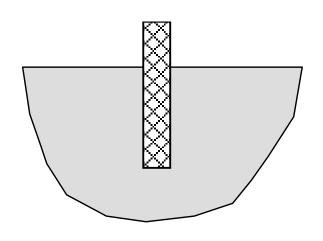
(Non-Coordinate Surface Boundary)

Line of Symmetry Boundary Conditions

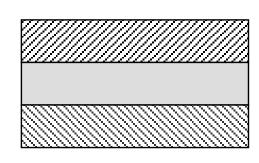


Line of symmetry boundary condition.

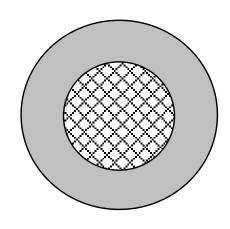
Interface Boundary Conditions



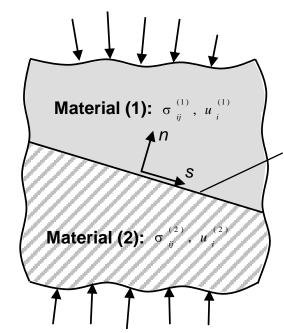




Layered Composite Plate



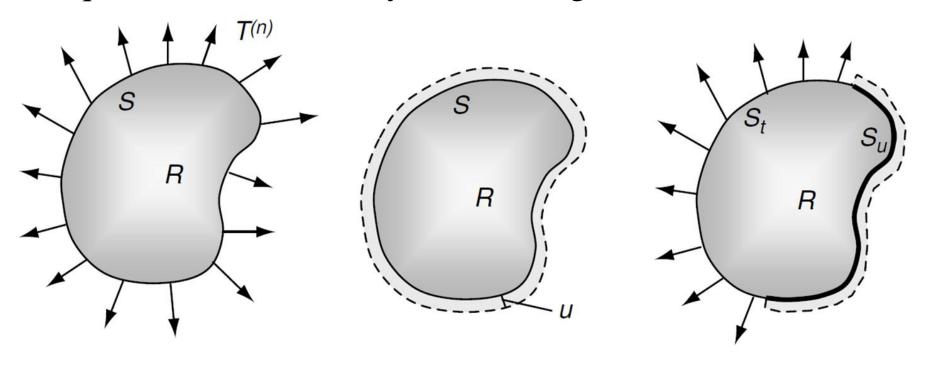
Composite Cylinder or Disk



Interface Conditions: Perfectly Bonded, Slip Interface, Etc.

Problem Classification

• Goal: determine the distribution of displacements, strains and stresses in the interior of an elastic body under equilibrium when body forces are given.



Traction Problem

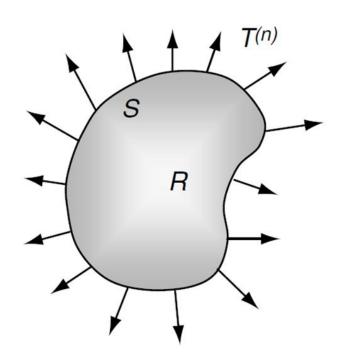
Displacement Problem

Mixed Problem

 $f\{u_{i}, \varepsilon_{ij}, \sigma_{ij}; \lambda, G, F_{i}\} = 0$

Simplification!

- Applicable to Traction Problem
- Boundary conditions are given in terms of the tractions or stress components
- Aiming to reformulate the field equations solely in terms of the stresses by eliminating the displacements and strains



• Using Hooke's law to rewrite the compatibility in terms of the stresses

$$\varepsilon_{ij} = \frac{1+v}{E} \sigma_{ij} - \frac{v}{E} \sigma_{kk} \delta_{ij} \quad \text{(Hooke's Law)}$$

$$\begin{cases} \varepsilon_{ij,kl} = \frac{1+v}{E} \sigma_{ij,kl} - \frac{v}{E} \sigma_{mm,kl} \delta_{ij}, & \varepsilon_{kl,ij} = \frac{1+v}{E} \sigma_{kl,ij} - \frac{v}{E} \sigma_{mm,ij} \delta_{kl} \\ \varepsilon_{ik,jl} = \frac{1+v}{E} \sigma_{ik,jl} - \frac{v}{E} \sigma_{mm,jl} \delta_{ik}, & \varepsilon_{jl,ik} = \frac{1+v}{E} \sigma_{jl,ik} - \frac{v}{E} \sigma_{mm,ik} \delta_{jl} \end{cases}$$

$$\varepsilon_{ij,kl} + \varepsilon_{kl,ij} - \varepsilon_{ik,jl} - \varepsilon_{jl,ik} = 0 \quad \text{(Strain compatibility)}$$

$$\Rightarrow \sigma_{ij,kl} + \sigma_{kl,ij} - \sigma_{ik,jl} - \sigma_{jl,ik} = \frac{v}{1+v} \left(\sigma_{mm,kl} \delta_{ij} + \sigma_{mm,ij} \delta_{kl} - \sigma_{mm,jl} \delta_{ik} - \sigma_{mm,ik} \delta_{jl} \right)$$

• Recall that, only 6 out of the 81 are meaningful

$$\boxed{k = l} \Rightarrow \boxed{\sigma_{ij,kk} + \frac{1}{1+v}\sigma_{kk,ij} = \frac{v}{1+v}\sigma_{mm,kk}\delta_{ij} + \sigma_{ik,jk} + \sigma_{jk,ik}}$$

• Further simplifying the compatibility using equilibrium equations

$$\sigma_{ik,k} = -F_{i} \Rightarrow \begin{cases} \sigma_{ik,jk} = -F_{i,j} \\ \sigma_{jk,ik} = -F_{j,i} \end{cases} \Rightarrow \sigma_{ij,kk} + \frac{1}{1+\nu} \sigma_{kk,ij} = \frac{\nu}{1+\nu} \sigma_{mm,kk} \delta_{ij} - (F_{i,j} + F_{j,i}) \end{cases}$$

$$i = j \Rightarrow \sigma_{ii,kk} = -\frac{1+\nu}{1-\nu} F_{i,i}$$

$$\Rightarrow \begin{bmatrix} \sigma_{ij,kk} + \frac{1}{1+\nu} \sigma_{kk,ij} = -\frac{\nu}{1-\nu} F_{k,k} \delta_{ij} - (F_{i,j} + F_{j,i}) \\ \nabla^{2} \sigma + \frac{1}{1+\nu} (\nabla \sigma_{kk}) \nabla = -\frac{\nu}{1-\nu} (\nabla \cdot \mathbf{F}) \mathbf{I} - (\mathbf{F} \nabla + \nabla \mathbf{F}) \end{bmatrix}$$

• Beltrami-Michell Compatibility Equations (Compatibility in terms of stress)

Scalar equations

$$\nabla^{2}\sigma_{x} + \frac{1}{1+\nu} \frac{\partial^{2}}{\partial x^{2}} (\sigma_{x} + \sigma_{y} + \sigma_{z}) = -\frac{\nu}{1-\nu} \left(\frac{\partial F_{x}}{\partial x} + \frac{\partial F_{y}}{\partial y} + \frac{\partial F_{z}}{\partial z} \right) - 2 \frac{\partial F_{x}}{\partial x},$$

$$\nabla^{2}\sigma_{y} + \frac{1}{1+\nu} \frac{\partial^{2}}{\partial y^{2}} (\sigma_{x} + \sigma_{y} + \sigma_{z}) = -\frac{\nu}{1-\nu} \left(\frac{\partial F_{x}}{\partial x} + \frac{\partial F_{y}}{\partial y} + \frac{\partial F_{z}}{\partial z} \right) - 2 \frac{\partial F_{y}}{\partial y},$$

$$\nabla^{2}\sigma_{z} + \frac{1}{1+\nu} \frac{\partial^{2}}{\partial z^{2}} (\sigma_{x} + \sigma_{y} + \sigma_{z}) = -\frac{\nu}{1-\nu} \left(\frac{\partial F_{x}}{\partial x} + \frac{\partial F_{y}}{\partial y} + \frac{\partial F_{z}}{\partial z} \right) - 2 \frac{\partial F_{z}}{\partial z},$$

$$\nabla^{2}\tau_{xy} + \frac{1}{1+\nu} \frac{\partial^{2}}{\partial x \partial y} (\sigma_{x} + \sigma_{y} + \sigma_{z}) = -\left(\frac{\partial F_{x}}{\partial y} + \frac{\partial F_{y}}{\partial x} \right),$$

$$\nabla^{2}\tau_{xz} + \frac{1}{1+\nu} \frac{\partial^{2}}{\partial x \partial z} (\sigma_{x} + \sigma_{y} + \sigma_{z}) = -\left(\frac{\partial F_{x}}{\partial z} + \frac{\partial F_{z}}{\partial x} \right),$$

$$\nabla^{2}\tau_{yz} + \frac{1}{1+\nu} \frac{\partial^{2}}{\partial x \partial z} (\sigma_{x} + \sigma_{y} + \sigma_{z}) = -\left(\frac{\partial F_{x}}{\partial z} + \frac{\partial F_{z}}{\partial x} \right),$$

• Only 3 are independent, i.e. • 3 equilibrium equations

$$\frac{\partial^{4}}{\partial x^{2} \partial y^{2}} \left(\sigma_{z} - v\left(\sigma_{x} + \sigma_{y}\right)\right) \\
= (1 + v) \frac{\partial^{3}}{\partial x \partial y \partial z} \left(-\frac{\partial \tau_{xy}}{\partial z} + \frac{\partial \tau_{yz}}{\partial x} + \frac{\partial \tau_{zx}}{\partial y}\right), \qquad \frac{\partial \sigma_{x}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + \frac{\partial \tau_{xz}}{\partial z} + F_{x} = 0,$$

$$\frac{\partial^{4}}{\partial y^{2} \partial z^{2}} \left(\sigma_{x} - v\left(\sigma_{y} + \sigma_{z}\right)\right) \qquad \frac{\partial \tau_{xy}}{\partial z} + \frac{\partial \sigma_{y}}{\partial z} + \frac{\partial \tau_{yz}}{\partial z} + F_{y} = 0,$$

$$= (1 + v) \frac{\partial^{3}}{\partial x \partial y \partial z} \left(-\frac{\partial \tau_{yz}}{\partial x} + \frac{\partial \tau_{zx}}{\partial y} + \frac{\partial \tau_{xy}}{\partial z}\right), \qquad \frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial z} + F_{z} = 0.$$

$$\frac{\partial^{4}}{\partial z^{2} \partial x^{2}} \left(\sigma_{y} - v\left(\sigma_{z} + \sigma_{x}\right)\right) \qquad \frac{\partial \tau_{xz}}{\partial z} + \frac{\partial \tau_{yz}}{\partial z} + \frac{\partial \sigma_{yz}}{\partial z} + F_{z} = 0.$$

$$= (1 + v) \frac{\partial^{3}}{\partial x \partial y \partial z} \left(-\frac{\partial \tau_{zx}}{\partial y} + \frac{\partial \tau_{xy}}{\partial z} + \frac{\partial \tau_{yz}}{\partial z}\right).$$

$$\frac{\partial x}{\partial x} + \frac{\partial y}{\partial z} + \frac{\partial \sigma_{y}}{\partial z} + \frac{\partial \sigma_{yz}}{\partial z} + F_{y} = 0,$$

$$\frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \sigma_{z}}{\partial z} + F_{z} = 0.$$

For 6 stress components

• The derivation of strains and displacements



Hooke's Law Strain-displacement relations

$$\left[\varepsilon_{ij} = \frac{1+\nu}{E}\sigma_{ij} - \frac{\nu}{E}\sigma_{kk}\delta_{ij}\right] \qquad \left[\varepsilon_{ij} = \frac{1}{2}\left(u_{i,j} + u_{j,i}\right)\right]$$

$$\varepsilon_{ij} = \frac{1}{2} \left(u_{i,j} + u_{j,i} \right)$$

Displacement Formulation

• Rewrite the Hooke's law in terms of displacements.

$$\left. \begin{array}{l} \varepsilon_{ij} = \frac{1}{2} \left(u_{i,j} + u_{j,i} \right); \varepsilon_{kk} = u_{k,k} \\ \\ \sigma_{ij} = \lambda \varepsilon_{kk} \delta_{ij} + 2 G \varepsilon_{ij} \end{array} \right\} \Rightarrow \boxed{\sigma_{ij} = \lambda u_{k,k} \delta_{ij} + G \left(u_{i,j} + u_{j,i} \right)}$$

Reformulate the equilibrium equations

$$\sigma_{ij,j} + F_i = 0$$

$$\Rightarrow \lambda u_{k,kj} \delta_{ij} + G \left(u_{i,jj} + u_{j,ij} \right) + F_i = 0$$

$$\Rightarrow G u_{i,kk} + (\lambda + G) u_{k,ki} + F_i = 0$$

$$\Rightarrow G \nabla^2 u + (\lambda + G) \nabla (\nabla \cdot u) + F = 0$$

Navier's/Lamé's Equations

Displacement Formulation

Scalar equations

$$G \nabla^2 u + (\lambda + G) \frac{\partial}{\partial x} \left[\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right] + F_x = 0,$$

$$G \nabla^2 v + (\lambda + G) \frac{\partial}{\partial y} \left[\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right] + F_y = 0,$$

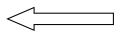
$$G \nabla^2 w + (\lambda + G) \frac{\partial}{\partial z} \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right) + F_z = 0.$$

• 3 equations for 3 displacements

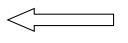
Displacement Formulation

The derivation of strains and stresses











Hooke's Law

Strain-displacement relations

$$\sigma = \lambda \operatorname{Tr}(\varepsilon) I + 2G \varepsilon$$

$$\sigma_{ij} = \lambda \varepsilon_{kk} \delta_{ij} + 2G \varepsilon_{ij}$$

$$\lambda = \frac{E v}{(1 + v)(1 - 2v)},$$

$$G = \frac{E}{2(1+v)}$$

$$\varepsilon = \frac{1}{2} \left(\mathbf{u} \nabla + \nabla \mathbf{u} \right)$$

$$\varepsilon_{ij} = \frac{1}{2} \left(u_{i,j} + u_{j,i} \right)$$

Summary of Formulations

General Field Equation System

(18 Equations, 15 Unknowns:)

$$f\{u_{i}, \varepsilon_{ij}, \sigma_{ij}; \lambda, G, F_{i}\} = 0$$

$$\varepsilon_{ij} = \frac{1}{2} \left(u_{i,j} + u_{j,i} \right);$$

$$\sigma_{ij} = \lambda \varepsilon_{kk} \delta_{ij} + 2 G \varepsilon_{ij};$$

$$\sigma_{ij,j} + F_i = 0;$$

$$\varepsilon_{ij,kl} + \varepsilon_{kl,ij} - \varepsilon_{ik,jl} - \varepsilon_{jl,ik} = 0.$$

Stress Formulation

(6 Equations, 6 Unknowns)

$$f\{\sigma_{ij}; v, F_i\} = 0$$

$$\sigma_{ij,j} + F_i = 0;$$

$$\sigma_{ij,kk} + \frac{1}{1+\nu} \sigma_{kk,ij}$$

$$= -\frac{\nu}{1-\nu} F_{k,k} \delta_{ij} - (F_{i,j} + F_{j,i}).$$

Displacement Formulation

(3 Equations, 3 Unknowns)

$$f\{u_i, \lambda, G, F_i\} = 0$$

$$\left|G u_{i,kk} + (\lambda + G) u_{k,ki} + F_i = 0\right|$$

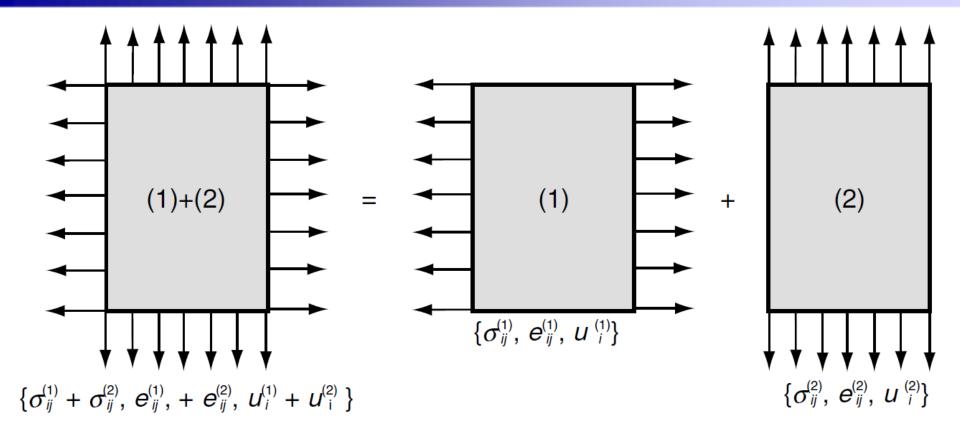
FEM codes!

Principle of Superposition

- Superposition applies to any problem that is governed by linear equations.
- Under the assumption of small deformations and linear elastic constitutive behavior, all elasticity field equations are linear.
- The usual displacement and traction boundary conditions are also linear.

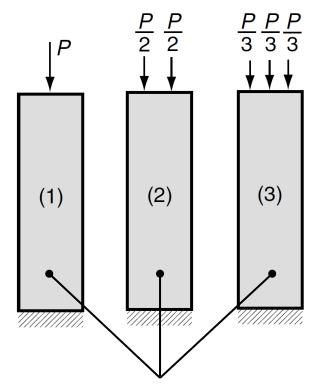
Principle of Superposition: For a given problem domain, if the state $\{\sigma_{ij}^{(1)}, e_{ij}^{(1)}, u_i^{(1)}\}$ is a solution to the fundamental elasticity equations with prescribed body forces $F_i^{(1)}$ and surface tractions $T_i^{(1)}$, and the state $\{\sigma_{ij}^{(2)}, e_{ij}^{(2)}, u_i^{(2)}\}$ is a solution to the fundamental equations with prescribed body forces $F_i^{(2)}$ and surface tractions $T_i^{(2)}$, then the state $\{\sigma_{ij}^{(1)} + \sigma_{ij}^{(2)}, e_{ij}^{(1)} + e_{ij}^{(2)}, u_i^{(1)} + u_i^{(2)}\}$ will be a solution to the problem with body forces $F_i^{(1)} + F_i^{(2)}$ and surface tractions $T_i^{(1)} + T_i^{(2)}$.

Principle of Superposition

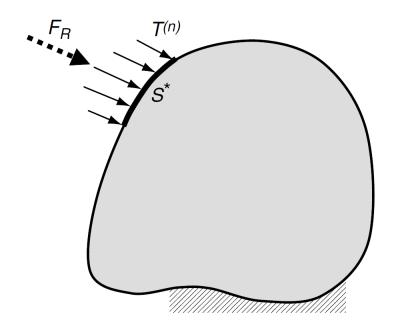


- Must have the same geometry and support type.
- Allow us to solve many more problems by using simpler basic cases whose solutions are already known.

Saint-Venant's Principle



Stresses approximately the same



Boundary loading $T^{(n)}$ would produce detailed and characteristic effects only in the vicinity of S^* . Away from S^* the stresses would generally depend more on the *resultant* F_R of the tractions rather than on the exact distribution

Saint-Venant's Principle

Saint-Venant's Principle: The stress, strain, and displacement fields caused by two different statically equivalent force distributions on parts of the body far away from the loading points are approximately the same.

• Allow us to solve a simpler statically equivalent problem with simpler boundary conditions to get a solution (approximate) for the original, more complicated problem.

Solution Strategies – Direct Method

- Seeking the solution of field equations by direct integration.
- Boundary conditions are satisfied exactly.
- Method normally encounters significant mathematical difficulties thus limiting its application to problems with simple geometry.

Solution Strategies – Inverse Method

- Displacements or stresses are selected that satisfy field equations.
- Finding/guessing a solution to the governing equations
- Trying to find a problem whose boundary conditions can be satisfied by the solution
- It is sometimes difficult to construct solutions to a specific problem of practical interest.

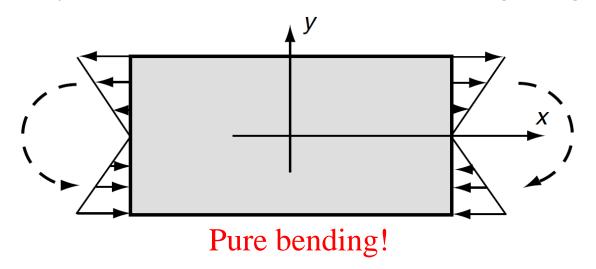
Sample Problem – Pure Bending of a Beam

☐ Consider the case of an elasticity problem under zero body forces with the following stress field:

$$\sigma_{x} = Ay$$
, $\sigma_{y} = \sigma_{z} = \tau_{xy} = \tau_{xz} = \tau_{yz} = 0$.

The stress field satisfies the equations of equilibrium and compatibility, and thus the field is a solution to an elasticity problem. What problem would be solved by such a field?

[Hint] Consider some trial domains and investigate the nature of the boundary conditions that would occur using the given stress field.



Solution Strategies – Semi-Inverse Method

- Part of displacement and/or stress field is specified, while the other remaining portion is determined by the fundamental field equations and the boundary conditions.
- Making an educated guess for part of the solution; then using the governing equations and/or boundary conditions to determine the rest of the solution.
- The usefulness of this approach is greatly enhanced by employing Saint-Venant's principle, whereby a complicated boundary condition can be replaced by a simpler statically equivalent distribution.

Sample Problem – Free Torsion of a Noncircular Shaft

■ Based on the torsion problem, We propose the following displacement field:

$$u = -\alpha zy$$
, $v = \alpha zx$, $w = w(x, y)$.

• By using the strain-displacement relations and Hooke's law, the stress field yields

$$\sigma_{x} = \sigma_{y} = \sigma_{z} = \tau_{xy} = 0; \quad \tau_{xz} = G\left(\frac{\partial w}{\partial x} - \alpha y\right); \quad \tau_{yz} = G\left(\frac{\partial w}{\partial y} + \alpha x\right).$$

• Substituting these results in equilibrium equations produces

$$\frac{\partial w^2}{\partial x^2} + \frac{\partial w^2}{\partial y^2} = 0.$$
 (Navier's equation)

[Note]

By assuming part of the solution field, the remaining equations to be solved are greatly simplified.

A specific domain in the *x-y* plane along with appropriate boundary conditions is needed to complete the solution to a particular problem.

Mathematical Techniques

- Analytical Solution Procedures
- Power Series Method
- Fourier Method
- Integral Transform Method
- Complex Variable Method
- Approximate Solution Procedures
- Ritz Method
- Galerkin Method
- Numerical Solution Procedures
- Finite Difference Method (FDM)
- Finite Element Method (FEM)
- Boundary Element Method (BEM)

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